

# Exploring the QCD phase diagram with fluctuations of conserved charges

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## **BNL-Bi-CCNU Collaboration:**

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# Definitions

Expansion of the pressure:

$$\frac{p}{T^4} = \sum_{i,j,k=0}^{\infty} \frac{1}{i!j!k!} \chi_{ijk,0}^{BQS} \left(\frac{\mu_B}{T}\right)^i \left(\frac{\mu_Q}{T}\right)^j \left(\frac{\mu_S}{T}\right)^k$$

$X = B, Q, S$ : conserved charges

**Lattice**

$$\chi_n^X = \left. \frac{\partial^n [p/T^4]}{\partial (\mu_X/T)^n} \right|_{\mu_X=0}$$

generalized susceptibilities

⇒ only at  $\mu_X = 0$ !

**Experiment**

$$\begin{aligned} VT^3 \chi_2^X &= \langle (\delta N_X)^2 \rangle \\ VT^3 \chi_4^X &= \langle (\delta N_X)^4 \rangle - 3 \langle (\delta N_X)^2 \rangle^2 \\ VT^3 \chi_6^X &= \langle (\delta N_X)^6 \rangle \\ &\quad - 15 \langle (\delta N_X)^4 \rangle \langle (\delta N_X)^2 \rangle \\ &\quad + 30 \langle (\delta N_X)^2 \rangle^3 \end{aligned}$$

cumulants of net-charge fluctuations

$$\delta N_X \equiv N_X - \langle N_X \rangle$$

⇒ only at freeze-out  $(\mu_f(\sqrt{s}), T_f(\sqrt{s}))$ !

# Motivations

## Explore the QCD phase diagram

- Analyze higher order cumulants and test universal scaling behavior
  - Explore the QCD phase diagram

This talk

## Analyze freeze-out conditions

- Match various cumulant ratios of measured fluctuations to QCD
  - determine freeze-out parameter

BNL-Bielefeld, PRL 109 (2012) 192302.

## Identify the relevant degrees of freedom

- Compare (lattice) QCD fluctuations to various hadronic/quasiparticle models:
  - deconfinement vs. chiral transition (melting of open strange/charm hadrons)
  - evidence for experimentally not yet observed hadrons



See talk by  
H.-T. Ding

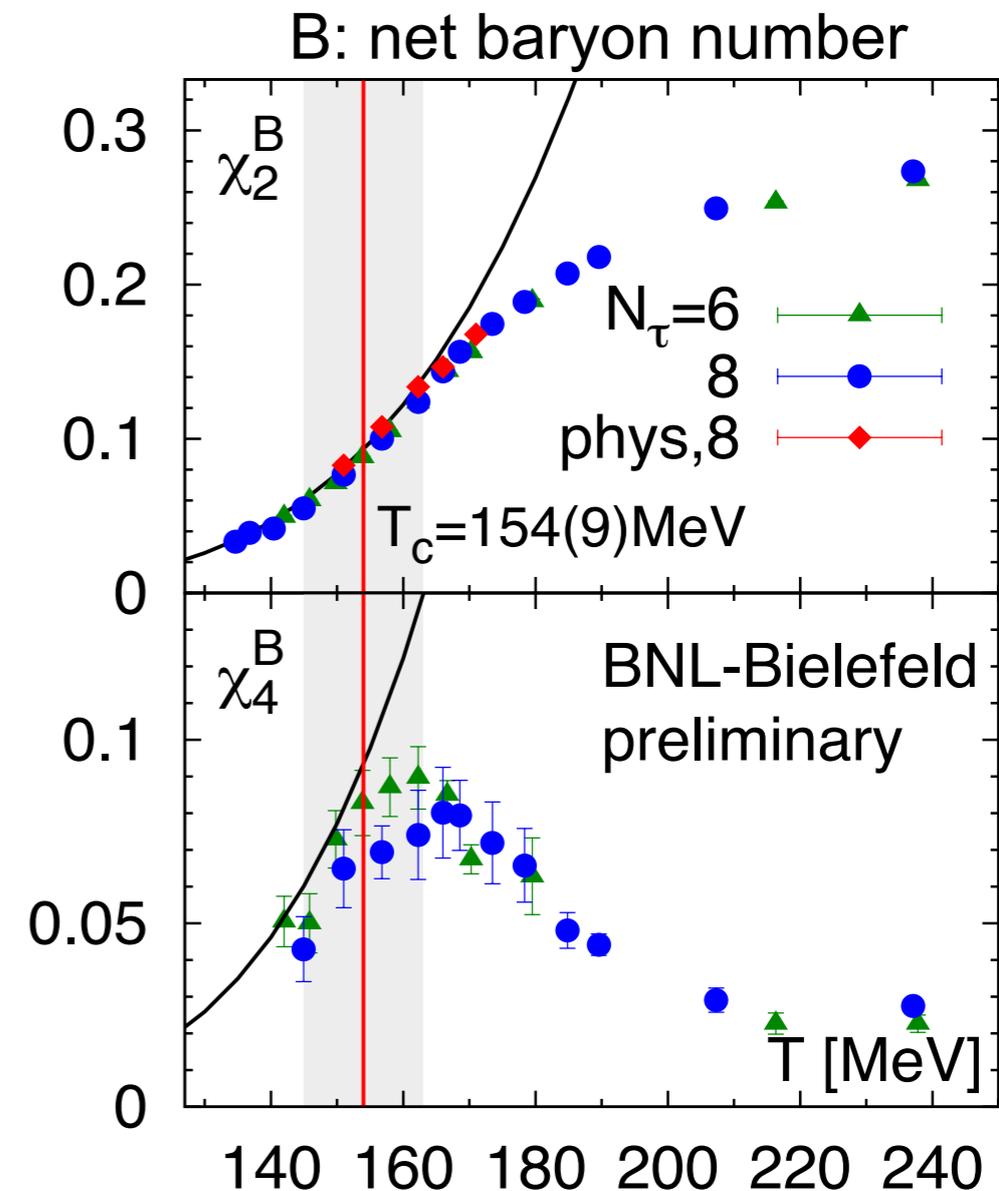
BNL-Bielefeld, PRL 111 (2013) 082301;

BNL-Bielefeld-CCNU, arXiv: 1404.4043, arXiv: 1404.6511.

# The lattice setup

## Lattice parameters:

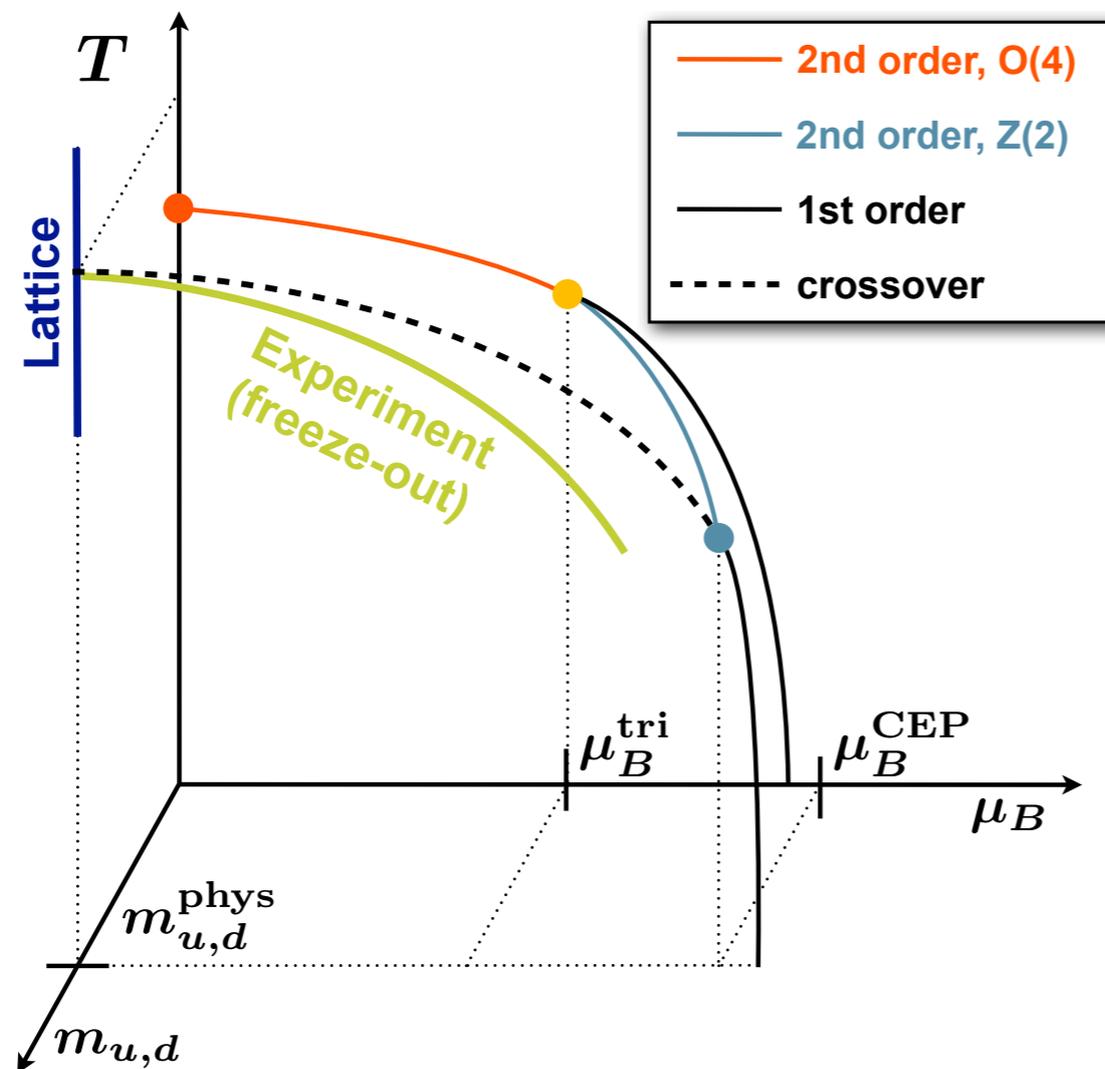
- (2+1)-flavor of highly improved staggered fermions (HISQ-action)
- a set of different lattice spacings ( $N_\tau = 6, 8, 12$ )
- two different pion masses:  $m_\pi = 140, 160 \text{ MeV}$
- high statistics:  $(10 - 16) \times 10^3$  configurations



⇒ statistical and systematical errors are under control

⇒ In general: find good agreement with HRG model for  $T < 155 \text{ MeV}$

# QCD critical behavior



assume scaling hypothesis for the free energy:

$$f = f_s(t, h) + \text{regular}$$

with

$$\lambda f_s(t, h) = f_s(\lambda^{a_t} t, \lambda^{a_h} h)$$

$\Rightarrow$  in the chiral limit ( $h=0$ ):

$$f \sim A_{\pm} |t|^{2-\alpha} + \text{regular}$$

critical exponent:

$\alpha$

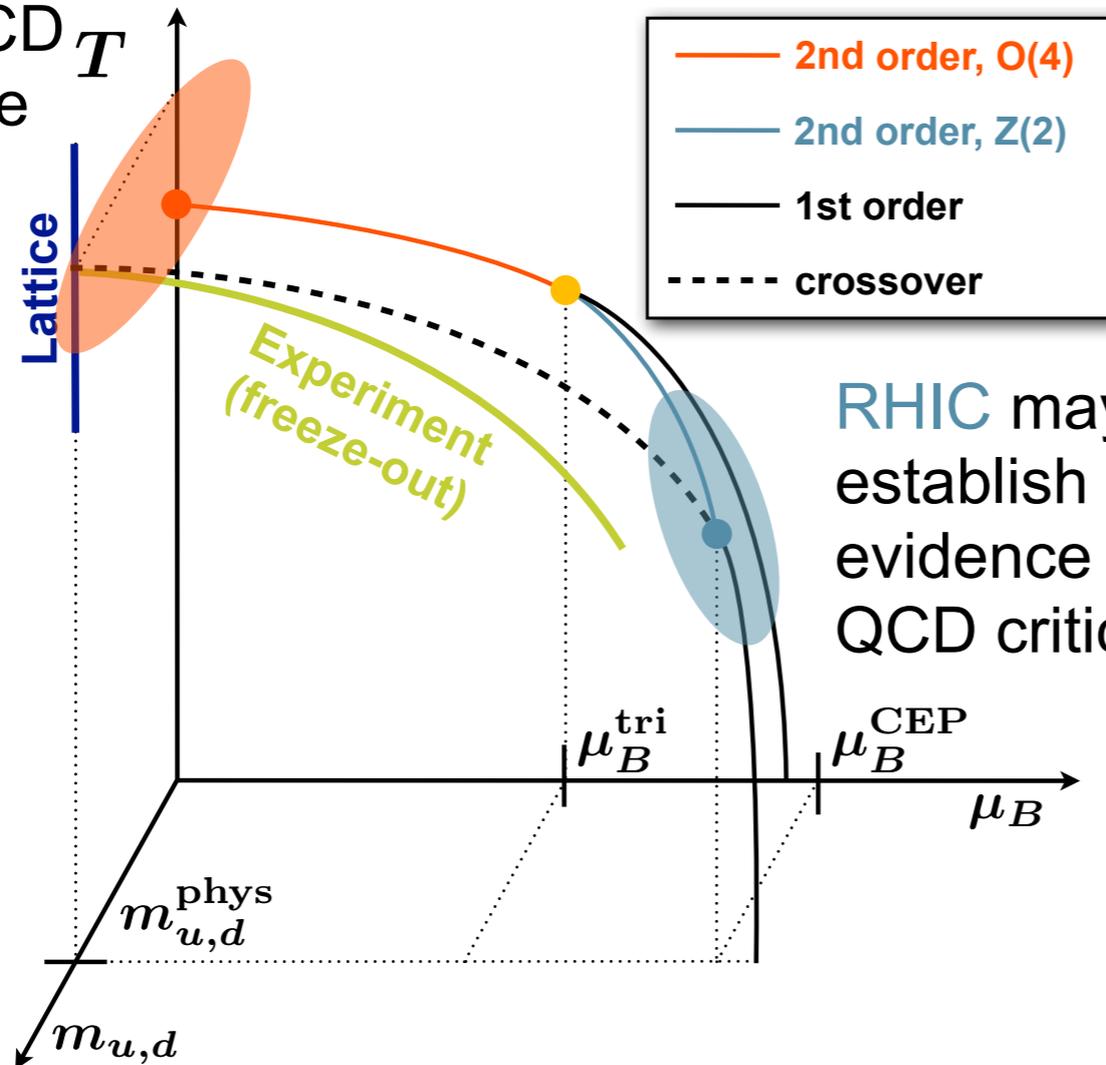
O(4) -0.213

Z(2) +0.107

$\Rightarrow$  at  $\mu_B = 0$ , 4th order cumulants develop a cusp, 6th order cumulants diverge

# QCD critical behavior

LHC may establish contact with the QCD chiral phase transition



RHIC may establish evidence for a QCD critical point

⇒ analyze universal scaling behavior

assume scaling hypothesis for the free energy:

$$f = f_s(t, h) + \text{regular}$$

with

$$\lambda f_s(t, h) = f_s(\lambda^{a_t} t, \lambda^{a_h} h)$$

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critical exponent:

	$\alpha$
O(4)	-0.213
Z(2)	+0.107

⇒ at  $\mu_B = 0$ , 4th order cumulants develop a cusp, 6th order cumulants diverge

# QCD critical behavior

matching scaling fields to QCD at  $\mu_B = 0$ :

$$h = \frac{m_q}{h_0} \quad t = \frac{1}{t_0} \left( \left( \frac{T - T_c}{T_c} \right) + \kappa \left( \frac{\mu_B}{T} \right)^2 \right)$$

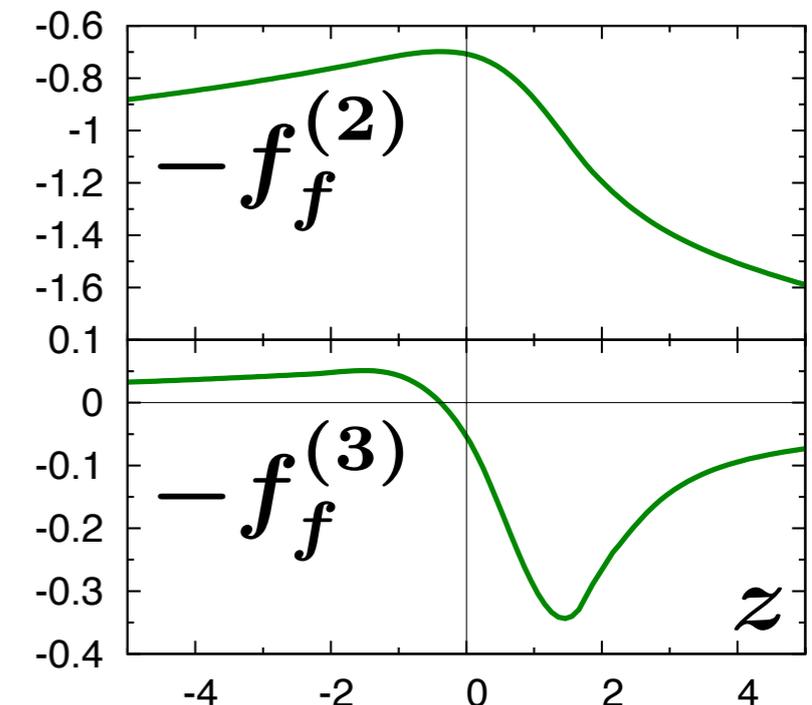
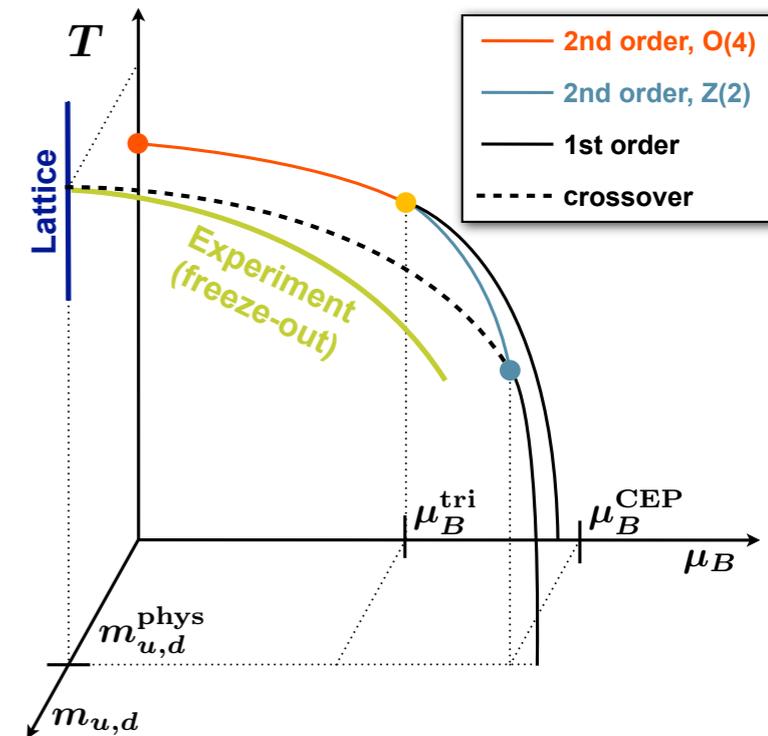
controlled by non-universal normalization constants  $t_0, h_0, \kappa$

$$\frac{p}{T^4} = -h^{(2-\alpha)/\beta\delta} \underbrace{f_f(t/h^{1/\beta\delta})}_{\text{(universal scaling function)}} - \underbrace{f_r(V, T, \vec{\mu})}_{\text{(regular part)}}$$

$f_f$  is function of scaling variable:  $z = t/h^{1/\beta\delta}$

$\Rightarrow$  critical behavior of cumulants (at  $\mu_B = 0$ ):

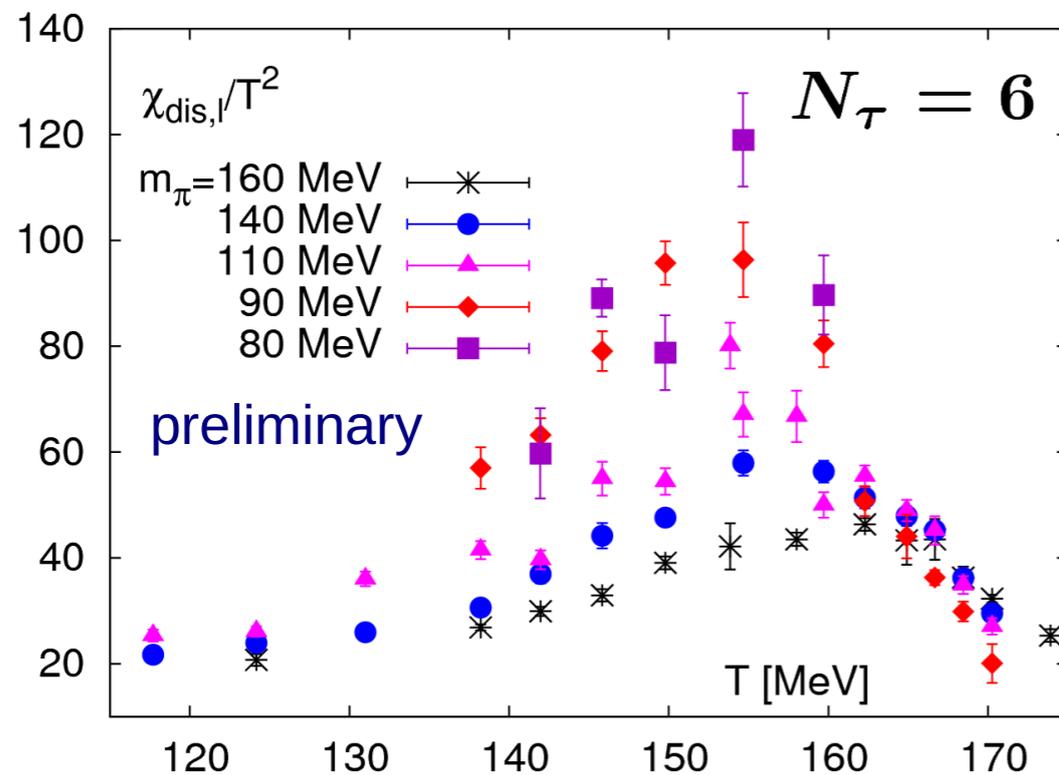
$$\chi_B^{(n)} \sim m_q^{(2-\alpha-n/2)/\beta\delta} f_f^{(n/2)}(t/h^{1/\beta\delta})$$



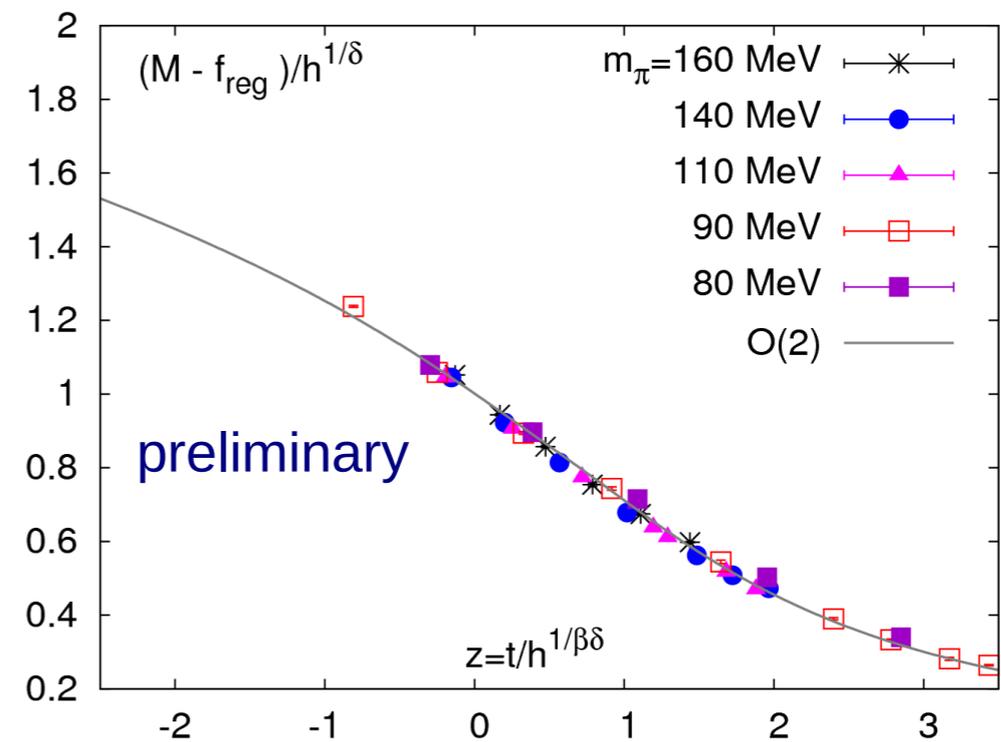
Karsch, Engels, PRD 85 (2012) 094506

# QCD critical behavior

evidence for universal scaling behavior with HISQ from chiral condensate and chiral susceptibility (H.-T. Ding et al. Lattice 2013)



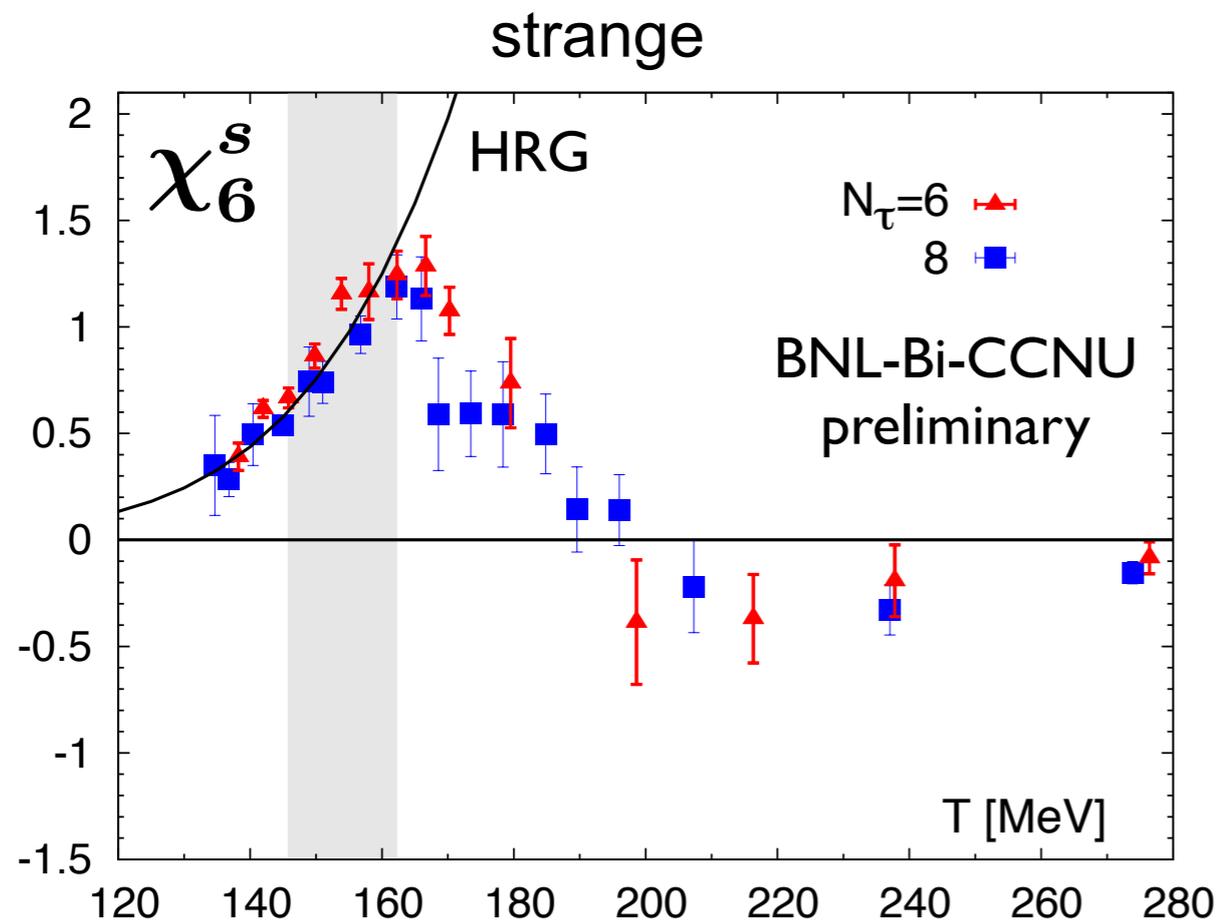
$$\chi_{dis}/T^2 \sim m_\pi^{2(1/\delta-1)}$$



$$M = m_s \langle \bar{\psi}\psi \rangle / T^4$$

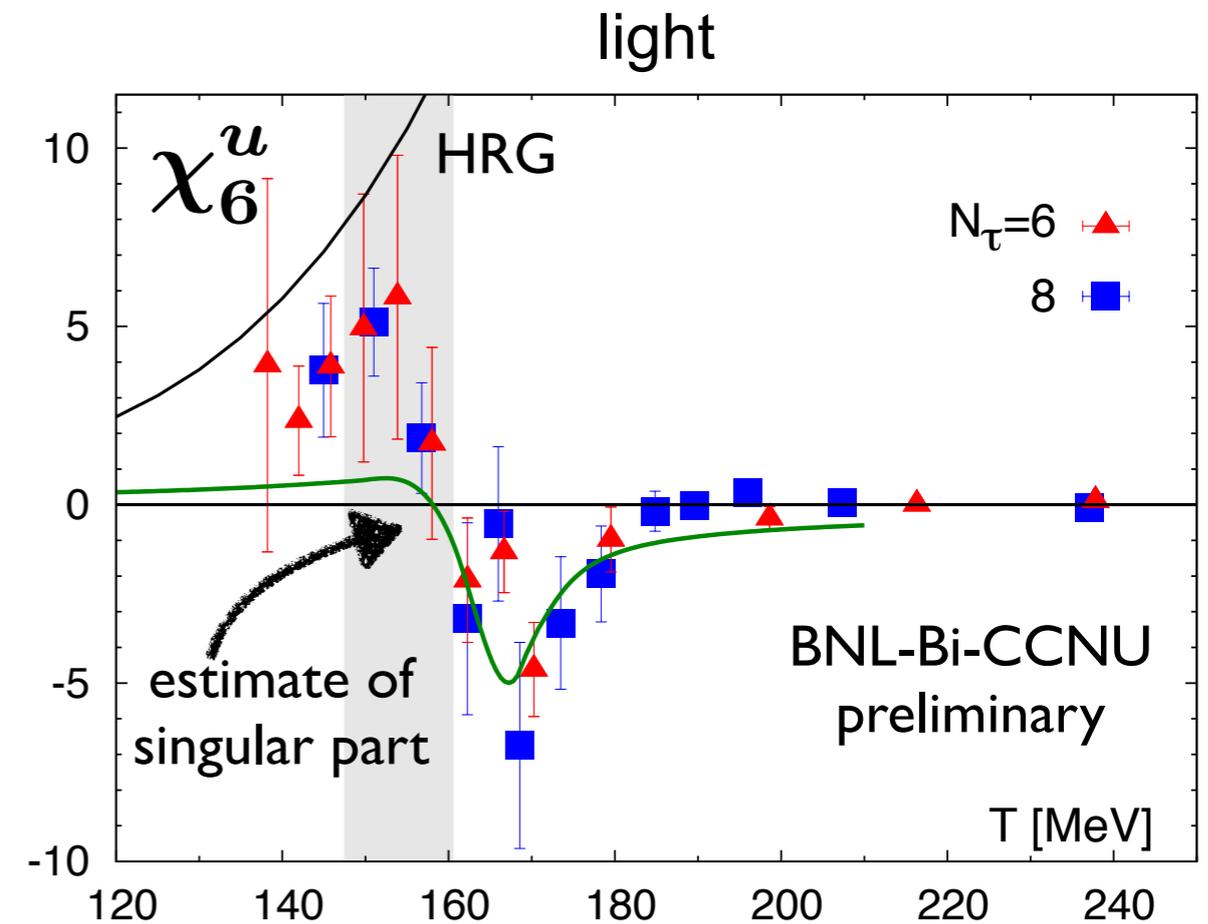
⇒ scaling region extends to physical pion mass

# QCD critical behavior



$\Rightarrow$  no evidence for typical O(4) singular structure

$\Rightarrow$  regular contribution dominates



$\Rightarrow$  clear evidence for typical O(4) singular structure

$\Rightarrow$  regular and singular contribution

# QCD critical behavior

some universal numbers:

width of the transition region (as seen by  $\chi_6^B$ ):

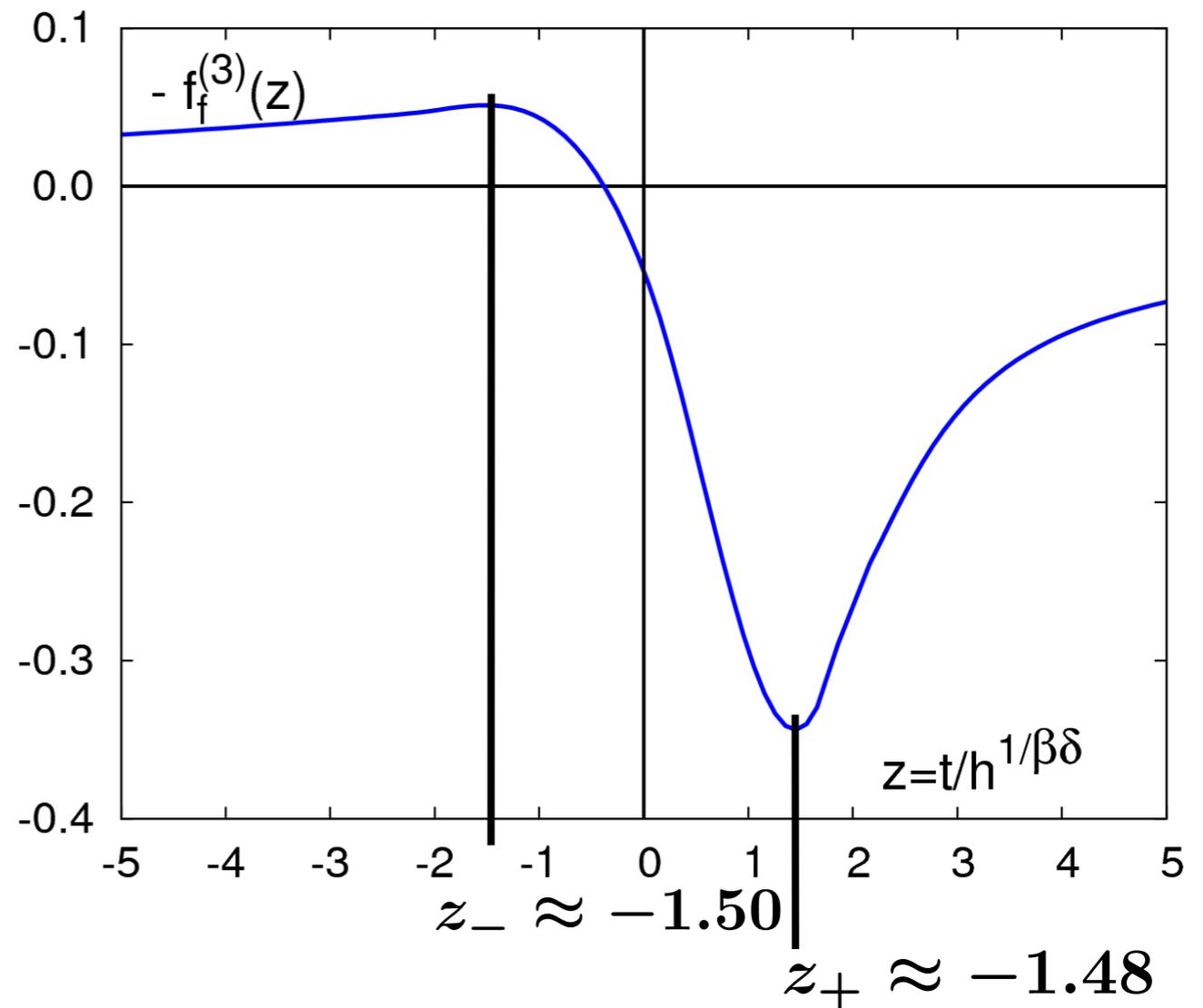
$$\Delta z = z_+ - z_- \approx 3$$

$\Rightarrow$

$$T_+ - T_- = \frac{1}{\Delta z} \frac{t_0 T_c}{h_0^{1/\beta\delta}} \left( \frac{m_l}{m_s} \right)^{1/\beta\delta}$$

at the physical point:

$$T_+ - T_- \approx 0.2 T_c$$



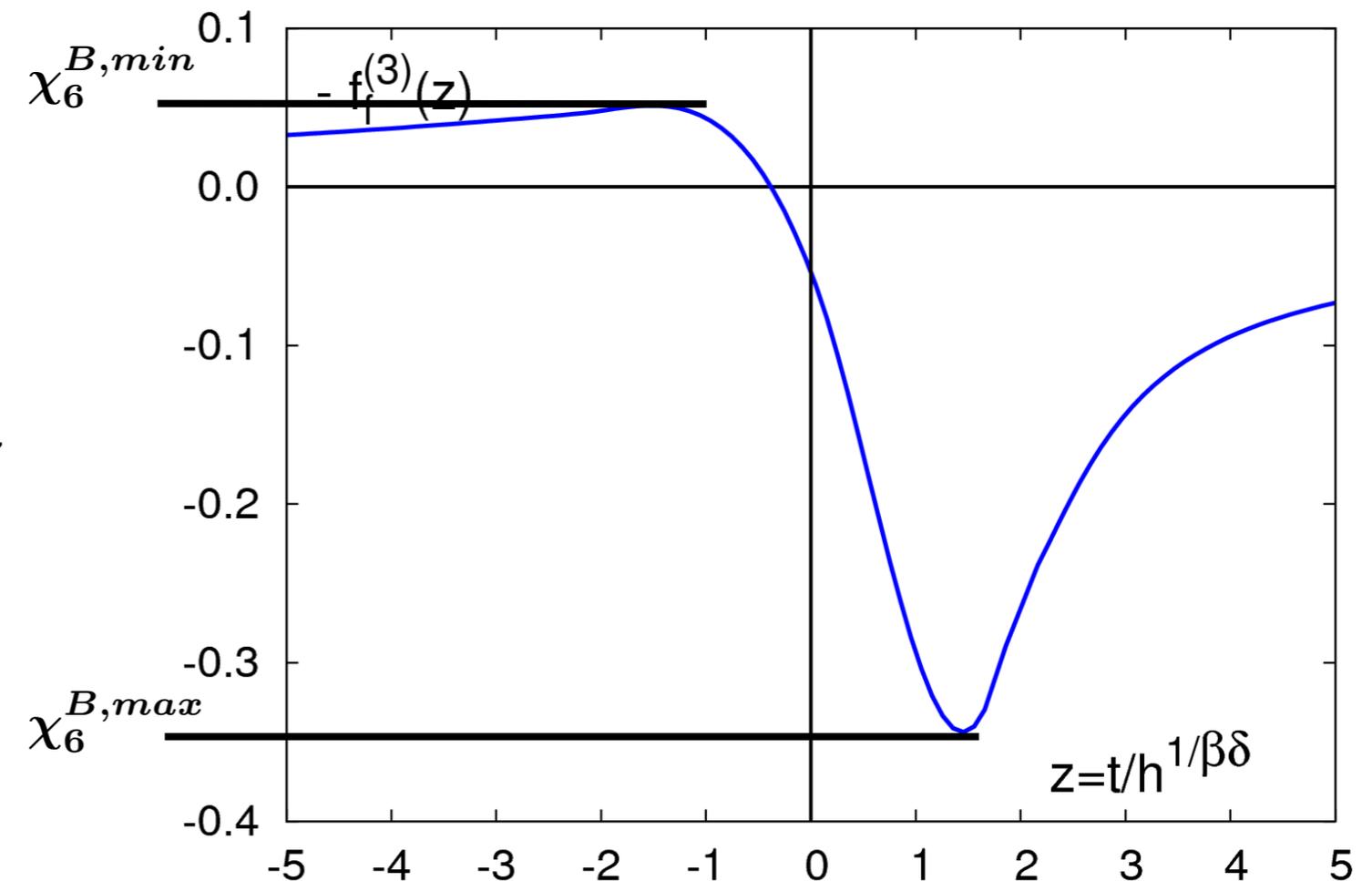
# QCD critical behavior

some universal numbers:

ratio of minimum to maximum:

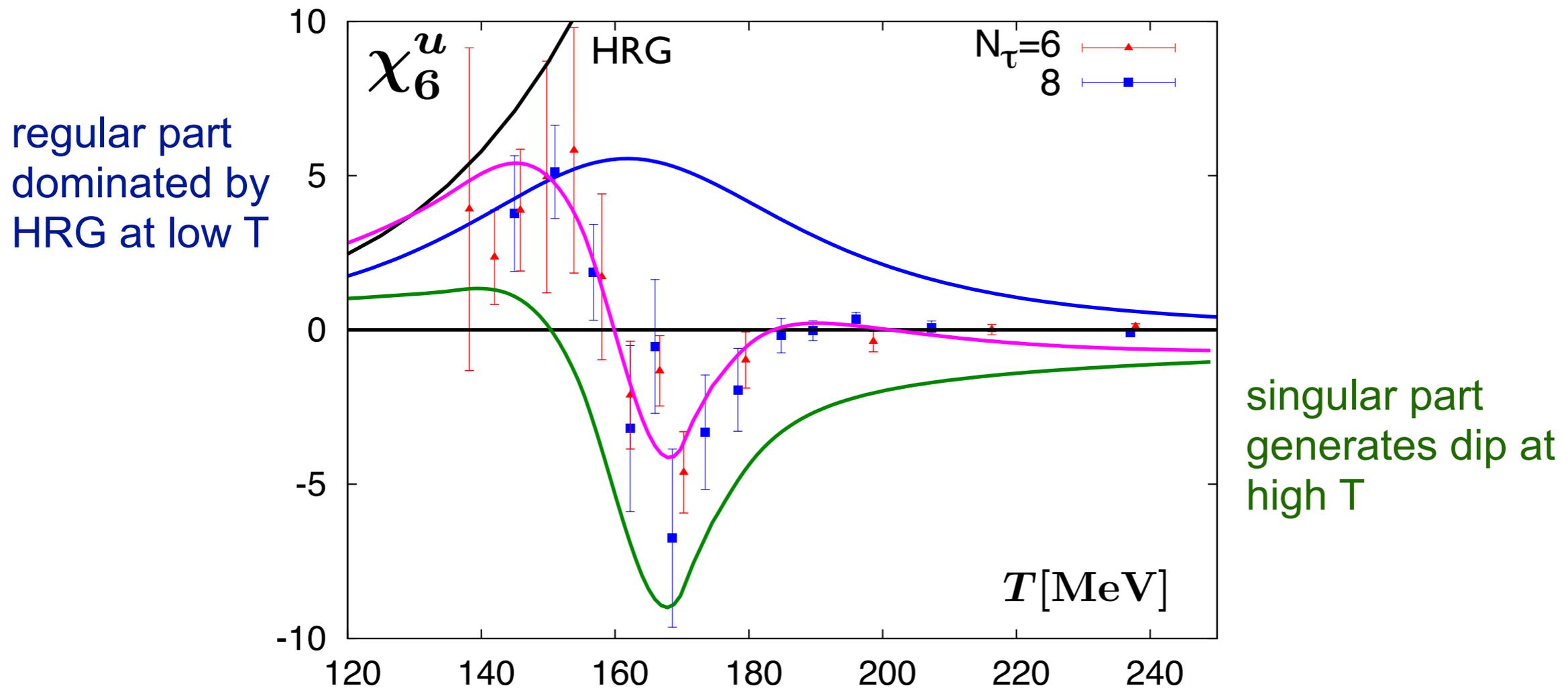
$$\chi_6^{B,min} / \chi_6^{B,max} \approx -6.7$$

⇒ depth of minimum at high T  
fixes maximal singular  
contribution at high T



# QCD critical behavior

on the interplay of regular and singular contributions (so far a guess and not a fit)

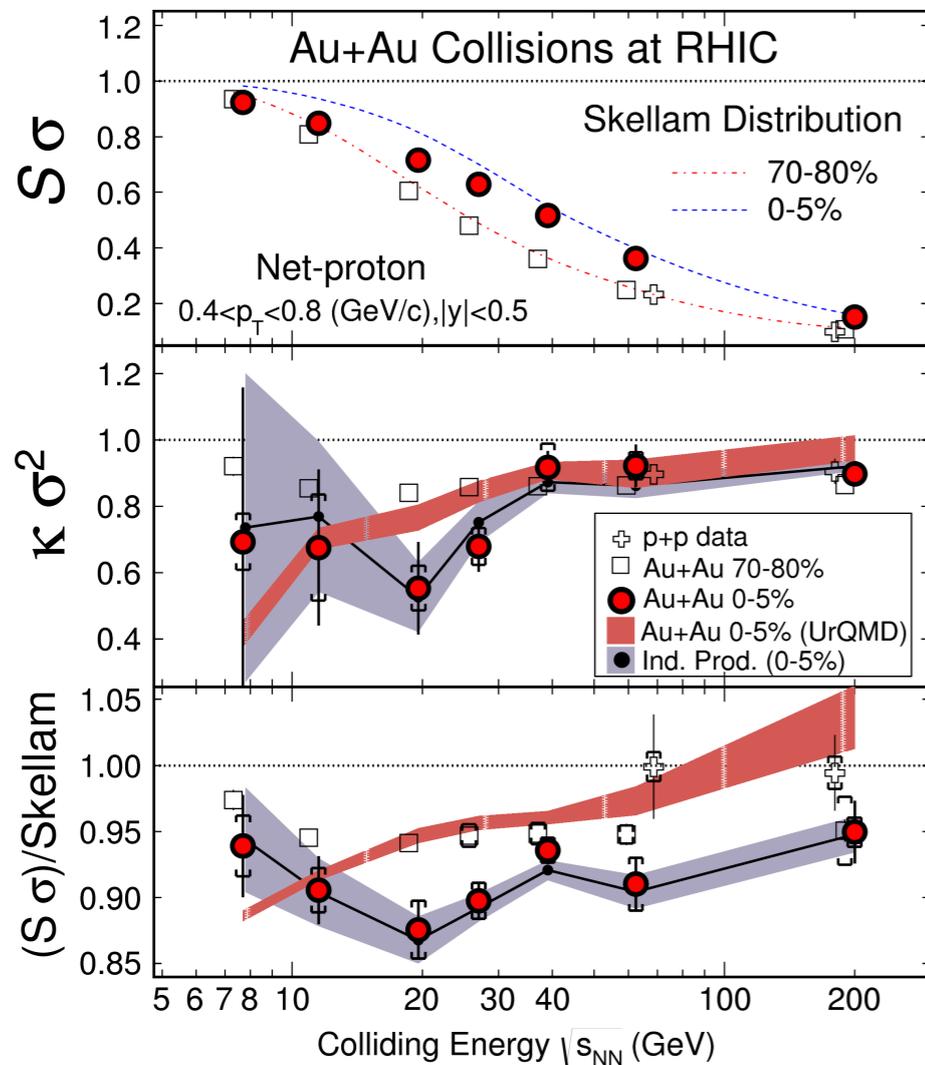


total = singular + regular

⇒ approximate agreement with HRG and sensibility to  $O(4)$  scaling are not in disagreement

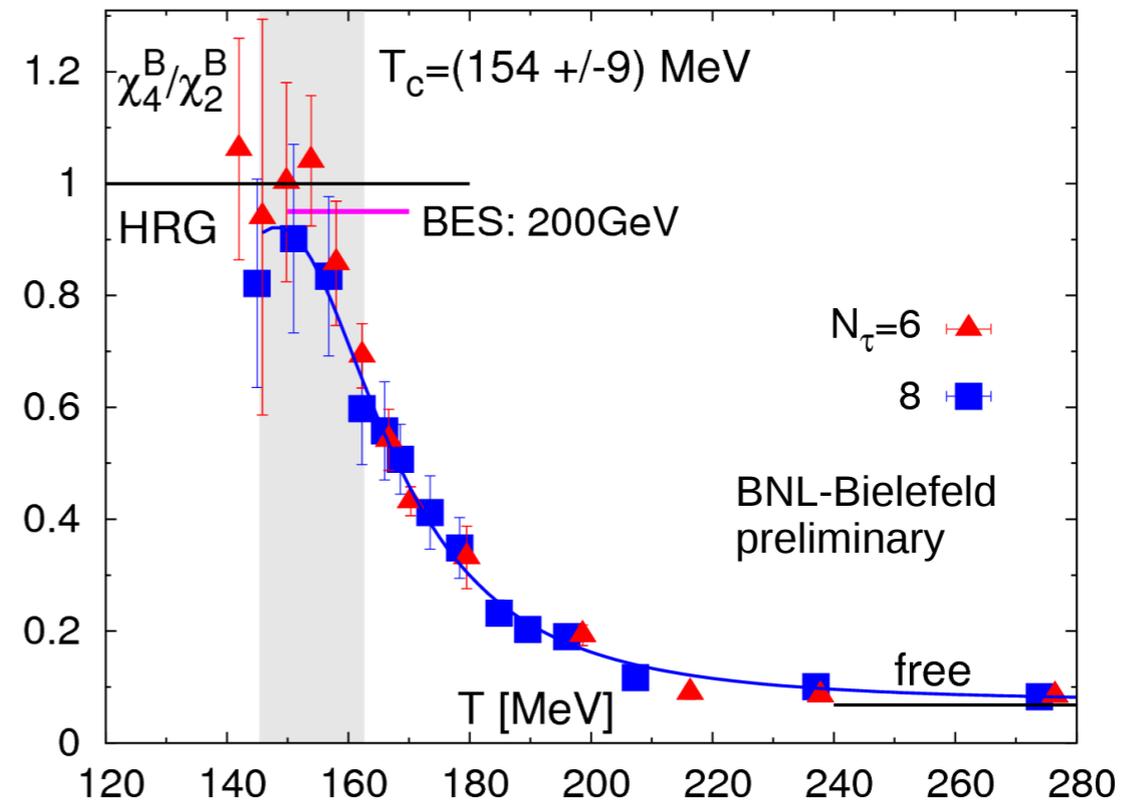
# Critical point search

proton number fluctuations



STAR, arXiv: 1309.5681

relative strength of the NLO correction to the pressure is controlled by  $\chi_4^B / \chi_2^B$



$$T < T_c : \quad 0.8 \leq \chi_4^B / \chi_2^B \leq 1.0$$

$$\Rightarrow \frac{\mathcal{O}(\mu_B^4) \text{ contribution to pressure}}{\mathcal{O}(\mu_B^2) \text{ contribution to pressure}} < 1$$

for  $\mu_B / T \lesssim 3.5$

# Critical point search

$$\frac{p}{T^4} = \frac{1}{VT^3} \ln Z(V, T, \mu_B) = \sum_{n, \text{even}} \frac{1}{n!} \chi_n^B \left( \frac{\mu_B}{T} \right)^n$$

⇒ consider radius of convergence

$$\left( \frac{\mu_B}{T} \right)_{crit} = \lim_{n \rightarrow \infty} \sqrt{\left| \frac{(n+2)! \chi_n^B}{n! \chi_{n+2}^B} \right|}$$

⇒ basic quantities

$$\chi_n^B / \chi_{n+2}^B$$

=1 for HRG

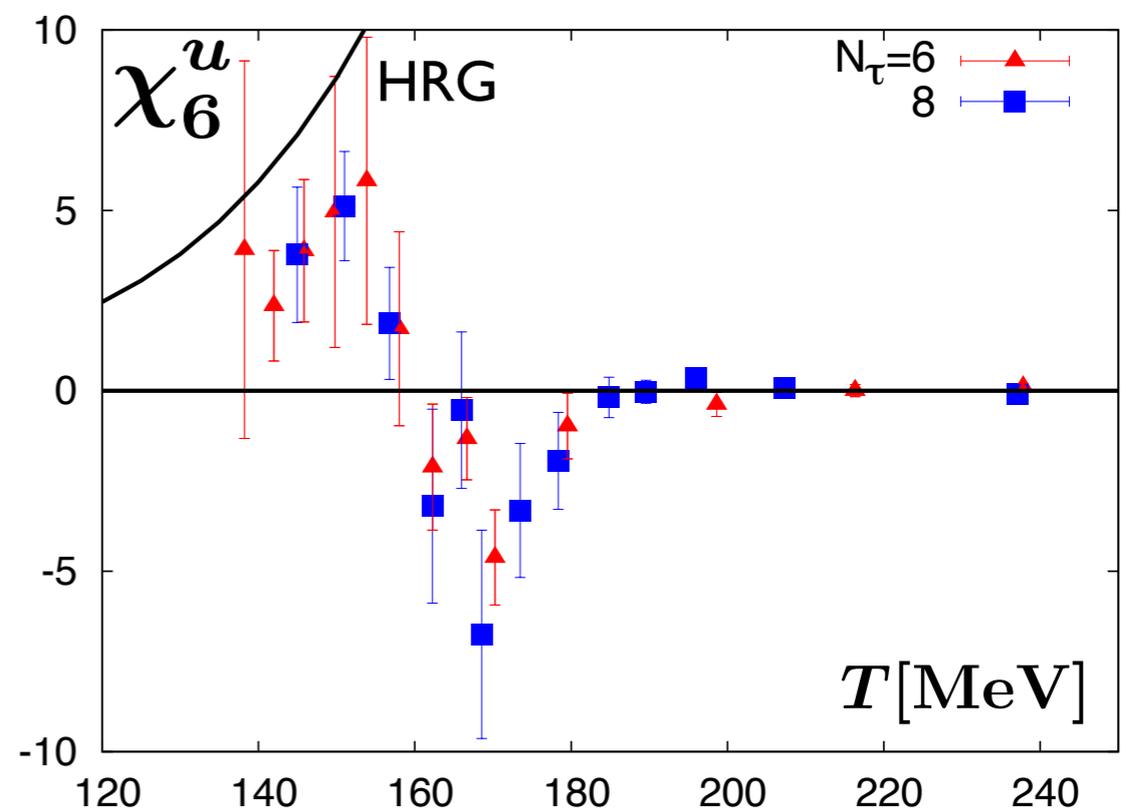
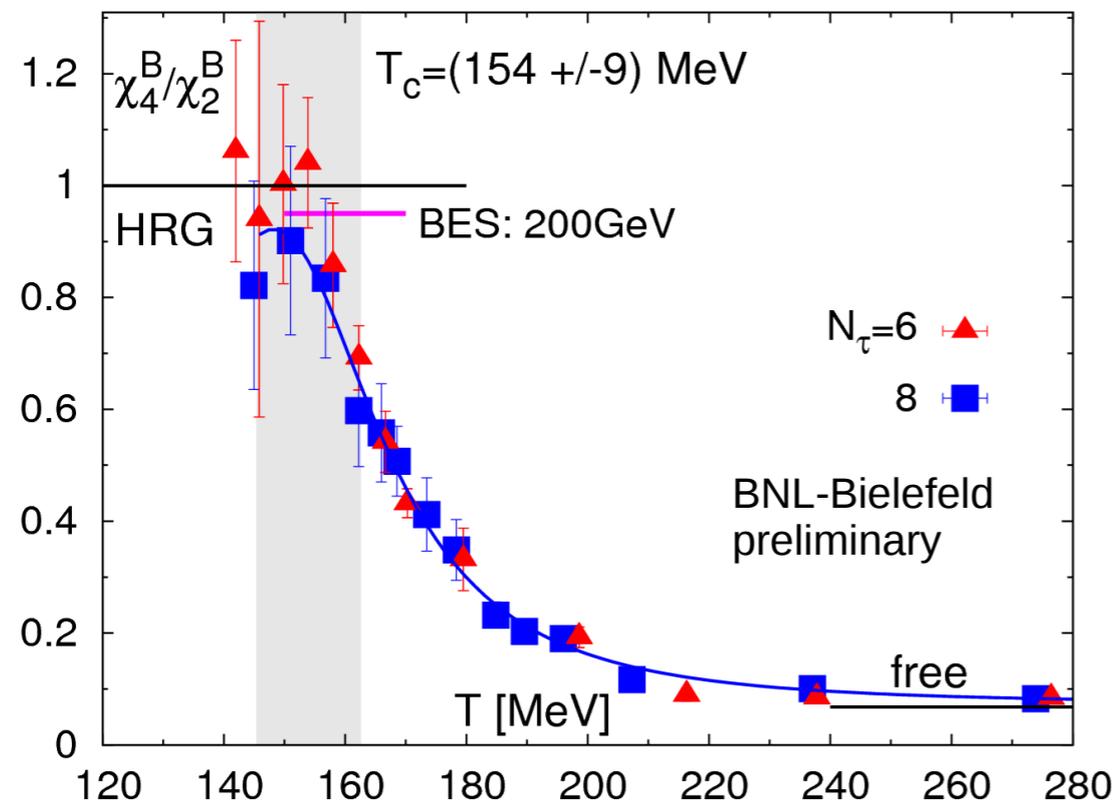
need to deviate from HRG like  $n^2$   
to obtain finite radius of convergence

⇒ singularity on the real axis **only if**

$$\chi_n^B > 0 \text{ for all } n > n_0$$

# Critical point search

we find so far no evidence for large enhancement over HRG for  $T < T_c$  ( $n \leq 6$ )



⇒ this suggests large  $\mu_B^{\text{crit}}$

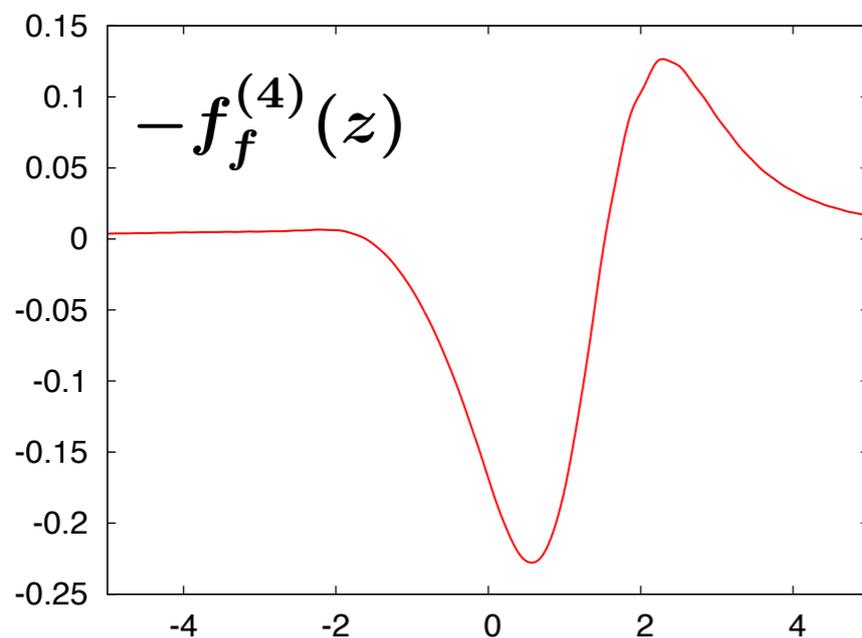
# Summary

- Approximate agreement with HRG model calculations at freeze-out and sensitivity to  $O(4)$  singular behavior are not inconsistent with each other
  - Higher order fluctuations need to deviate from HRG like  $n^2$  in order to obtain a finite radius of convergence
- 
- 6th order cumulants are sensitive to  $O(4)$  scaling but will pick up only a small singular contribution at low  $T$ . This favors estimates for the location of a critical end point at large baryon chemical potentials

# QCD critical behavior at $\mu_B > 0$

$$\begin{aligned} \mu_B > 0 : \chi_{4,\mu}^B = & -3(2\kappa_B t_0^{-1})^2 h^{-\alpha/\Delta} f_f^{(2)}(z) \\ & -6(2\kappa_B t_0^{-1})^3 (\hat{\mu}_B^c)^2 h^{-(1+\alpha)/\Delta} f_f^{(3)}(z) \\ & - (2\kappa_B t_0^{-1} \hat{\mu}_B^c)^4 h^{-(2+\alpha)/\Delta} f_f^{(4)}(z) + \text{regular} \end{aligned}$$

dominates in the chiral limit or if  $\hat{\mu}_B^c > 0 \gtrsim 1$



$\Rightarrow$  close to  $T_c$ :

$$\chi_4^B(\mu_B) < 0$$

B.Friman, FK, K.Redlich, V.Skokov,  
Eur. Phys. J. C71, 1694 (2011)

$\Rightarrow$  mapping of scaling variables non trivial

M. Stephanov, PRL 107 (2011) 052301